

Non-diffusive transport in pressure driven plasma turbulence

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Beyond the standard diffusive transport paradigm

- Experimental and theoretical evidence suggests that transport in fusion plasmas deviates from the diffusion paradigm:

$$\partial_t T = \partial_x [\chi(x,t) \partial_x T] + S(x,t)$$

- Examples:
- Fast propagation and non-local transport phenomena
 - Inward transport observed in off-axis fueling experiments
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- Our goal is to develop transport models that overcome the limitations of the diffusion paradigm.
 - The models incorporate in a unified way **non-locality**, **memory** effects, and **non-diffusive scaling**.
 - To motivate and test the models we consider transport in **pressure driven plasma turbulence**

Non-diffusive transport

Pressure gradient driven plasma turbulence

$$(\partial_t + \tilde{\mathbf{V}} \cdot \nabla) \nabla_{\perp}^2 \tilde{\Phi} = \frac{B_0}{m_i n_0 r_c} \frac{1}{r} \frac{\partial \tilde{p}}{\partial \theta} - \frac{1}{\eta m_i n_0 R_0} \nabla_{\parallel}^2 \tilde{\Phi} + \mu \nabla_{\parallel}^4 \tilde{\Phi}$$

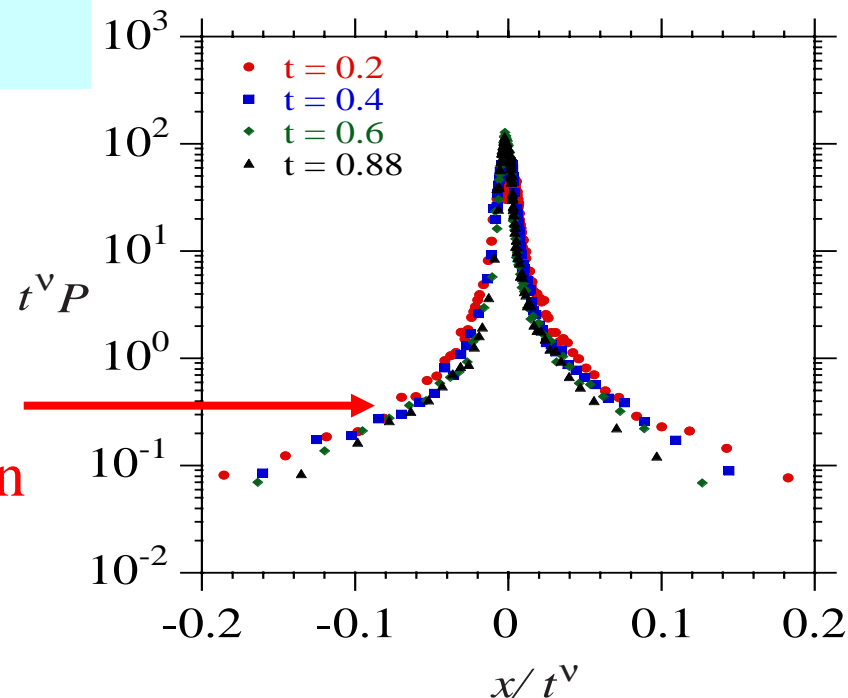
$$(\partial_t + \tilde{\mathbf{V}} \cdot \nabla) \tilde{p} = \frac{\partial \langle p \rangle}{\partial r} \frac{1}{r} \frac{\partial \tilde{\Phi}}{\partial \theta} + \chi_{\perp} \nabla_{\perp}^2 \tilde{p} + \chi_{\parallel} \nabla_{\parallel}^2 \tilde{p}$$

$$\frac{\partial \langle p \rangle}{\partial t} + \frac{1}{r} \frac{\partial}{\partial r} r \langle \tilde{V}_r \tilde{p} \rangle = S_0 + D \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \langle p \rangle}{\partial r} \right)$$

Tracers dynamics

$$\frac{d\vec{r}}{dt} = \tilde{\mathbf{V}} = \frac{1}{B^2} \nabla \tilde{\Phi} \times \vec{B}$$

Probability density function
of radial displacements



Super-diffusive
scaling

$$\langle \delta r^2 \rangle \sim t^{4/3}$$

Levy
distribution

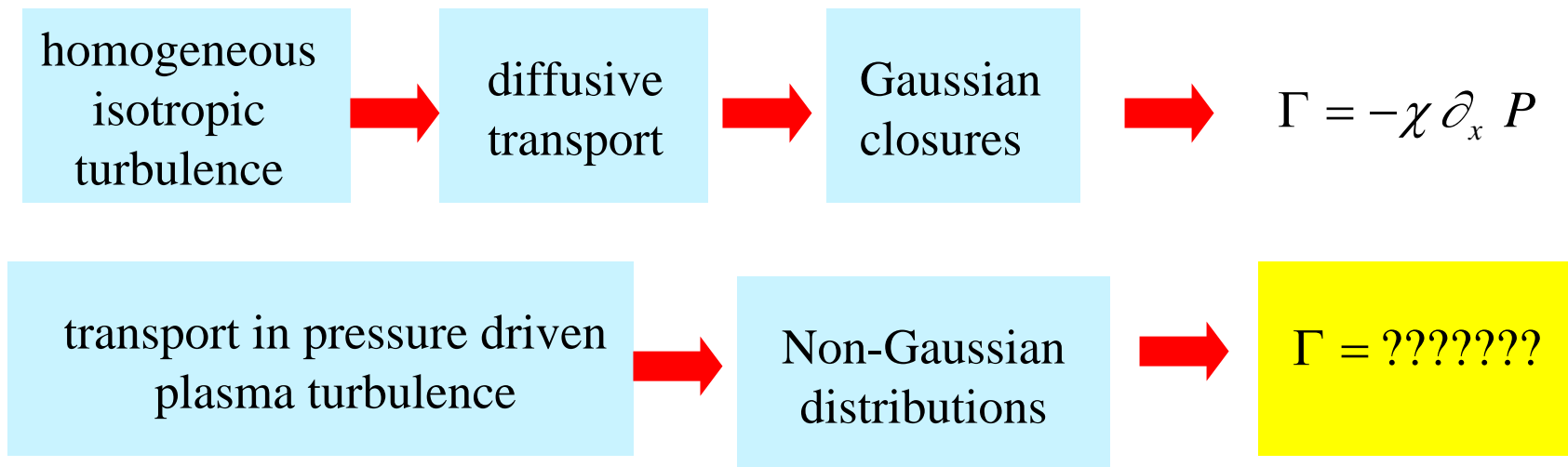
Towards an effective transport model for non-diffusive turbulent transport

- Individual tracers follow the turbulent field
- The distribution of tracers P evolves according to
- The idea is to construct a model that “encapsulates” the complexity of the turbulence field \tilde{V} in an effective flux Γ , and reproduces the observed pdf

$$\frac{d\vec{r}}{dt} = \tilde{V} = \frac{1}{B^2} \nabla \tilde{\Phi} \times \vec{B}$$

$$\frac{\partial P}{\partial t} + \tilde{V} \cdot \nabla P = 0$$

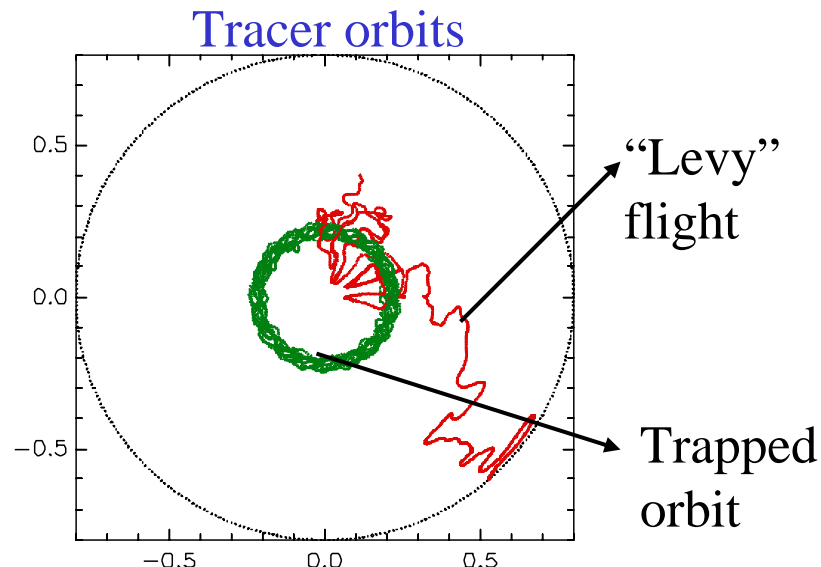
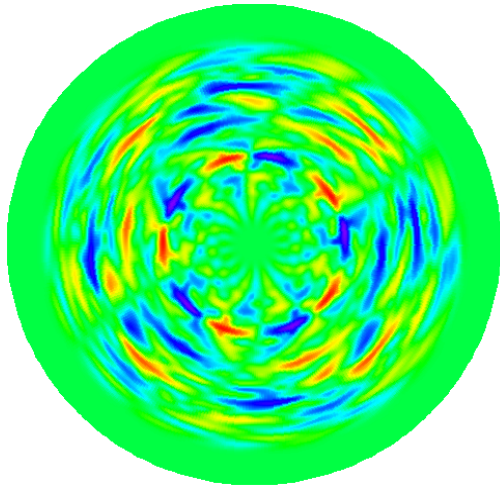
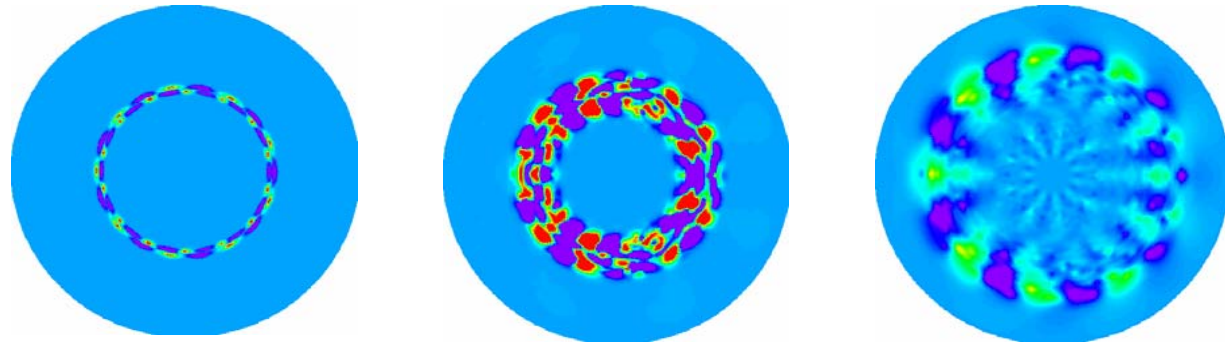
$$\frac{\partial P}{\partial t} = - \frac{\partial \Gamma}{\partial x}$$



What is the origin of non-diffusive transport?

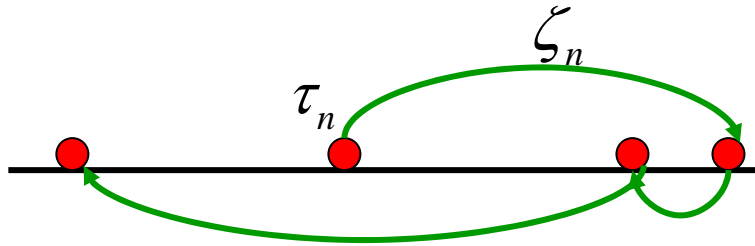
ExB flow velocity eddies induce large tracer trapping that leads to temporal non-locality, or **memory effects**

“Avalanche like” phenomena induce large tracer displacements that lead to spatial **non-locality**



The combination of tracer trapping and flights leads to **non-diffusive transport**

Continuous time random walk model



τ_n = waiting time $\psi(\tau)$ = waiting time pdf

ζ_n = jump $\lambda(\zeta)$ = jump size pdf

Master Equation

(Montroll-Weiss 1965)

$$P(x,t) = \underbrace{\delta(x) \int_t^\infty \psi(t') dt'}_{\text{Contribution from particles that have not moved during } (0,t)} + \underbrace{\int_0^t \psi(t-t') \left[\int_{-\infty}^\infty \lambda(x-x') P(x',t') dx' \right] dt'}_{\text{Contribution from particles located at } x' \text{ and jumping to } x \text{ during } (0,t)}$$

Contribution from particles that have not moved during $(0,t)$

Contribution from particles located at x' and jumping to x during $(0,t)$

No memory

$$\psi(\tau) \sim e^{-\mu\tau}$$

No long displacements

$$\lambda(\zeta) \sim e^{-\zeta^2/2\sigma^2}$$

$$\partial_t P = \partial_x [\chi \partial_x P] + S$$

Standard diffusion model

Proposed transport model

Long waiting times $\psi(\tau) \sim \tau^{-(\beta+1)}$

Long displacements
(Levy flights) $\lambda(\zeta) \sim \zeta^{-(\alpha+1)}$

$$\frac{\partial P}{\partial t} = - \frac{\partial}{\partial x} [w^- \Gamma_\ell + w^+ \Gamma_r]$$

Non-local effects
due to avalanches
causing Levy flights
modeled with **integral
operators in space.**

$$\Gamma_\ell = - \chi \frac{\partial}{\partial t} \int_0^t d\tau \int_a^x dy K(x-y; t-\tau) P(y, \tau)$$

$$K(x-y; t-\tau) = \frac{1}{(t-\tau)^{1-\beta} (x-y)^\alpha}$$

Non-Markovian,
memory effects
due to trapping in
eddies modeled
with **integral
operators in time.**

Equivalent formulation
using **fractional derivatives**

$$\frac{\partial^\beta P}{\partial t^\beta} = \chi \frac{\partial^\alpha P}{\partial x^\alpha}$$

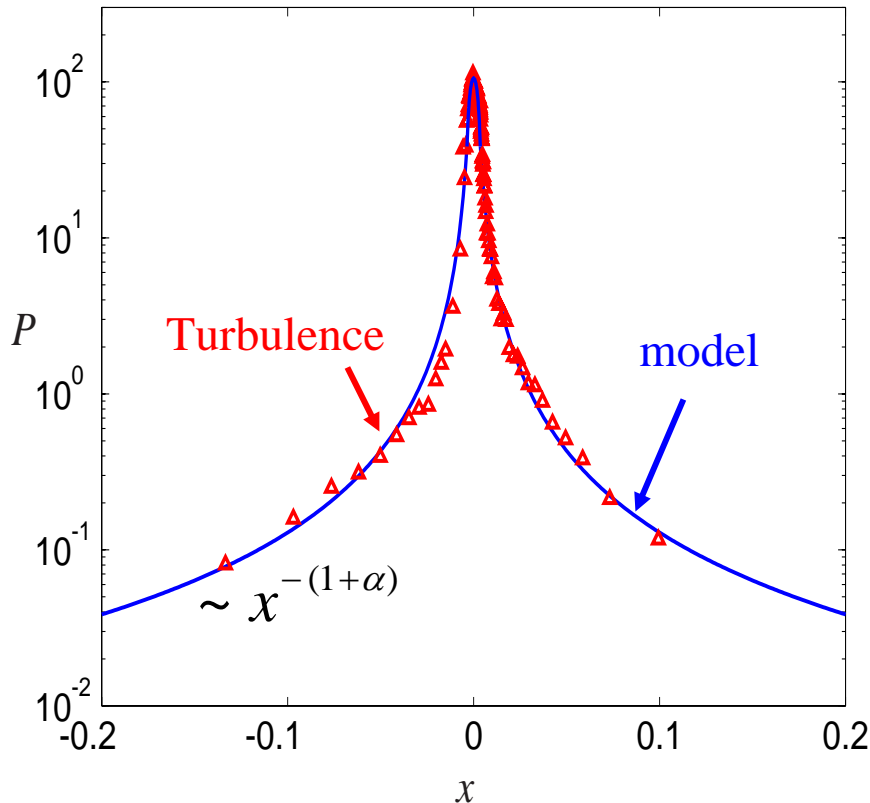
For pressure driven
plasma turbulence

$$\alpha = 3/4 \quad \beta = 1/2$$

Test of fractional transport model

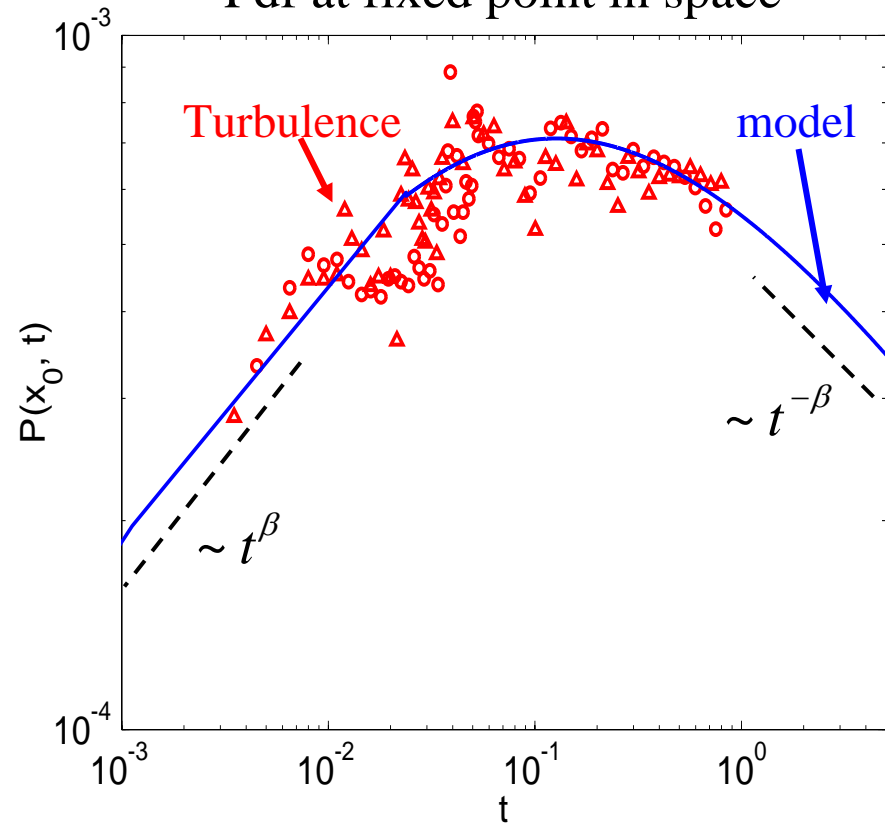
$$\frac{\partial^\beta P}{\partial t^\beta} = \chi \frac{\partial^\alpha P}{\partial x^\alpha}$$

Levy distribution at fixed time



Algebraic decay in space due to “Levy flights” implies that there is **no** characteristic transport scale

Pdf at fixed point in space



Algebraic scaling in time caused by **memory effects**

Super-diffusive scaling $\langle \delta r^2 \rangle \sim t^{2\nu}$

Diffusive scaling

$$\nu = 1/2$$

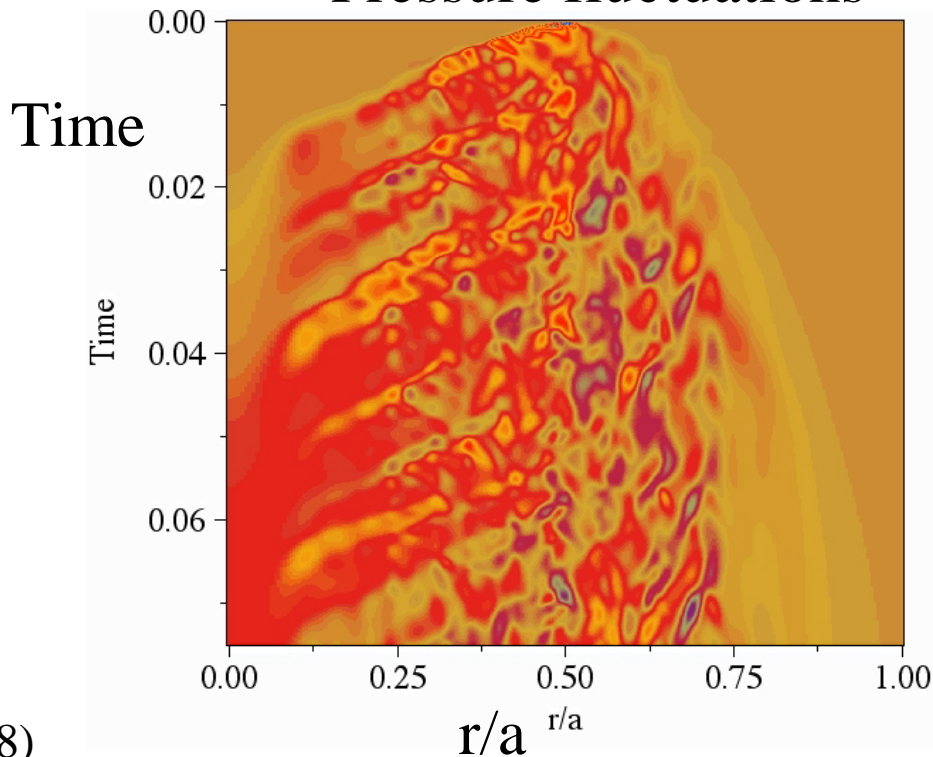
Tracers

$$\nu = 0.63$$

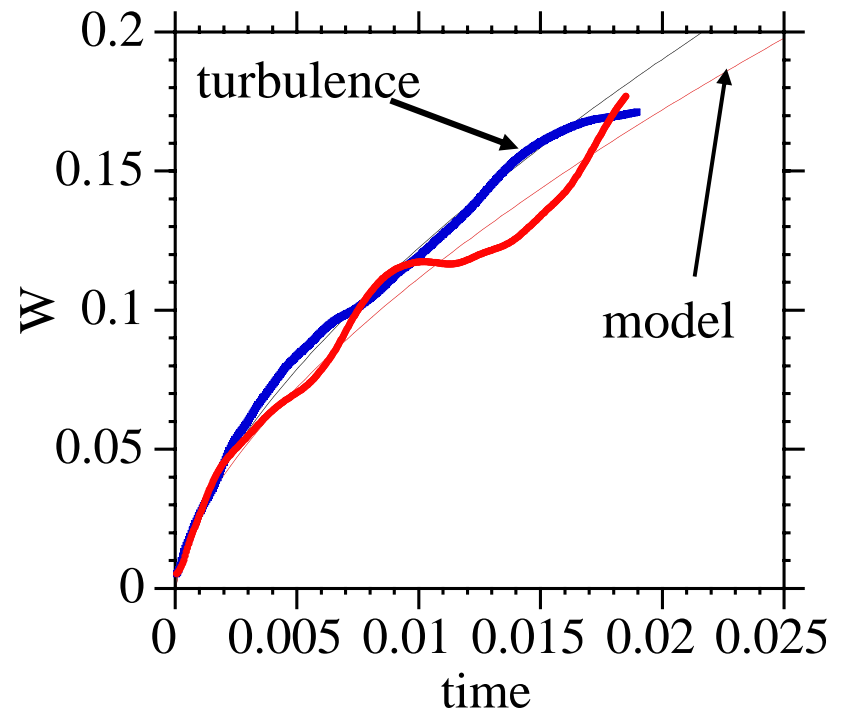
Model

$$\nu = \beta / \alpha = 2/3 = 0.666$$

Pressure fluctuations



$$W \sim \langle x^2 \rangle$$



Conclusions

- We presented numerical evidence of non-diffusive transport i.e., **super-diffusion** and **Levy distributions**, in plasma turbulence.
- We proposed a transport model that incorporates in a unified way space **non-locality**, **memory effects** and **anomalous diffusion scaling**.
- There is **quantitative agreement** between the model and the turbulence calculations.
- The model represents a first attempt to construct **effective transport operators** when the complexity of the turbulence invalidates the use of Gaussian closures.

$$\left[\partial_t + \tilde{V} \cdot \nabla \right] \Leftrightarrow \left[\partial_t^\beta - \chi \partial_x^\alpha \right] \quad \text{Fractional derivatives}$$


Further details: [del-Castillo-Negrete, et al., Phys. of Plasmas, 11, 3854, \(2004\).](#)