

Internal kink mode stability in the presence of ICRH driven fast ions populations

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Outline

Fast ions



$$\delta W_{MHD} \rightarrow \delta W_{MHD} + \delta W_{HOT}$$

Effect of ICRH driven populations on
Internal kink mode n=1, m=1

MHD energy
functional

Fast ions
energy
functional

Perturbative method (F. Porcelli et al. Phys. Plasmas, 1 (3), 470 (1994))	Quantitative results	Sawteeth
Variational method (R. White et al. Phys. Fluids B 2, 745 (1990))	Qualitative results	Sawteeth, Fishbones

Objective : Explain experimental observations in JET

Perturbative method

$$\longrightarrow \delta W = \delta W_{MHD} + \delta W_{HOT} \quad \gamma = \gamma_I + \gamma_{HOT}$$

↓ ↓
eigenfunction perturbation

Relevant equations, which include Finite Orbit Width effects,

$$\delta W_{HOT} = \delta W_{HOT}^{ad} + \delta W_{HOT}^{na}$$

(F. Porcelli *et al.*
Phys. Plasmas, 1
(3), 470 (1994))

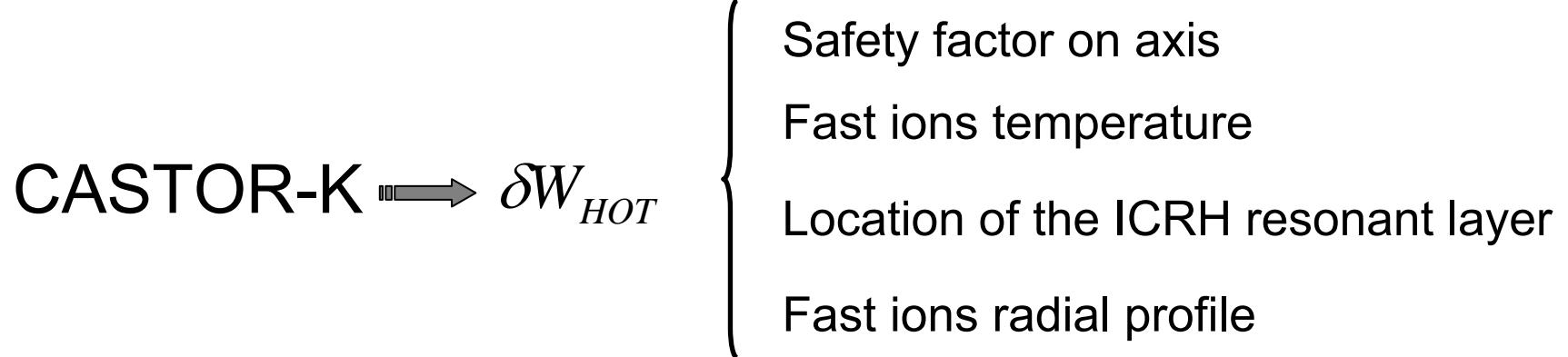
$$\delta W_{HOT}^{ad} = \frac{2\pi^2}{\Omega m^2} \sum_{\sigma} \int dE d\mu dP_{\phi} (\xi_{\perp} \cdot \nabla \psi) \frac{Ze}{c} \frac{\partial F}{\partial P_{\phi}}$$

Adiabatic part - destabilizing

Perturbative method

$$\delta W_{HOT}^{na} = -\frac{2\pi^2}{\Omega m^2} \sum_{\sigma} \int dE d\mu dP_{\varphi} \tau_b (\omega - n\omega_*) \frac{\partial F}{\partial E} \sum_{p=-\infty}^{\infty} \frac{|Y_p|^2}{\omega + n\langle \dot{\varphi} \rangle + p\omega_b}$$

Non-adiabatic part - stabilizing



$$\omega \ll \omega_D \quad \dots \rightarrow \gamma_{HOT} = -\frac{1}{2\gamma} \frac{\text{Re } \delta W_{HOT}}{E_k}$$

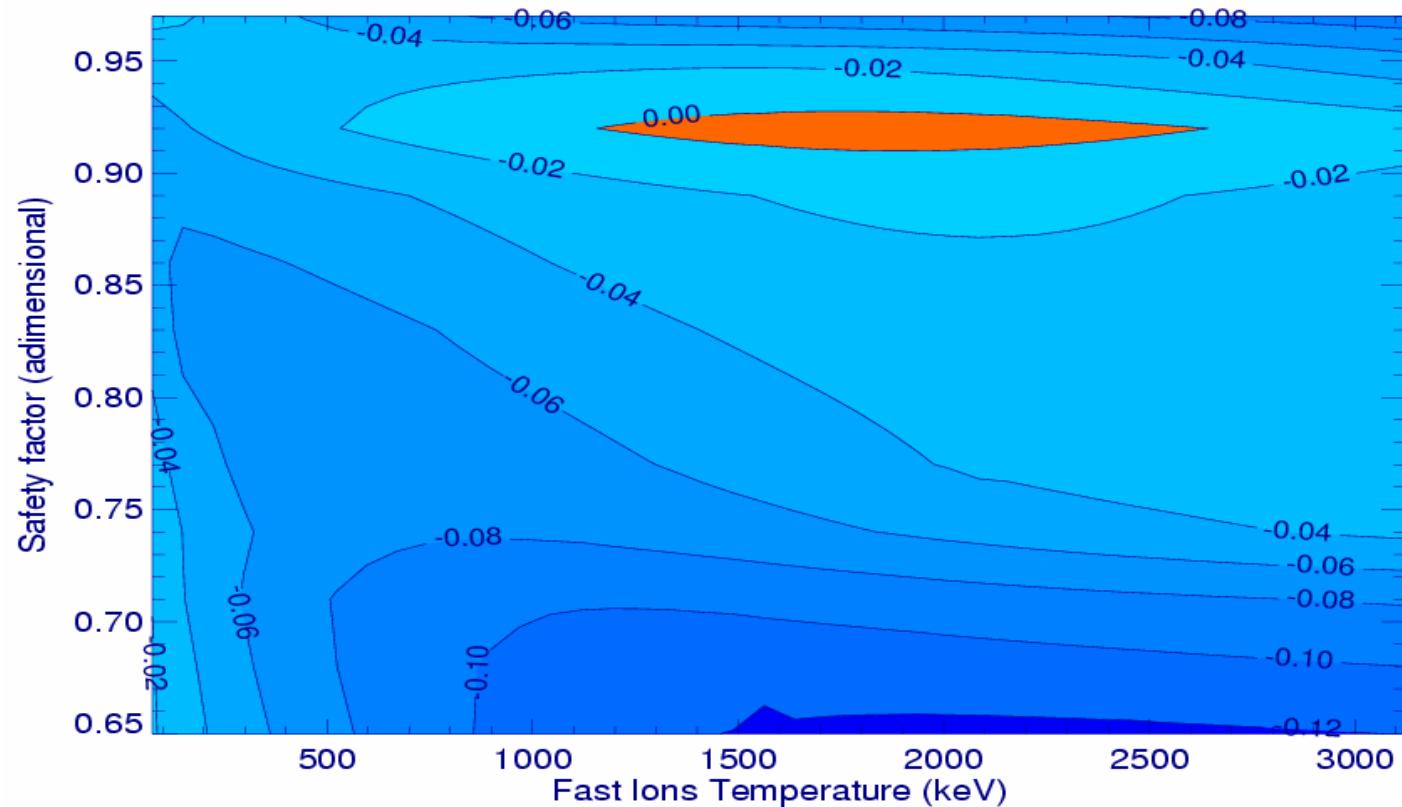
Sawteeth

Effect of the ICRH fast ions on sawteeth

Non-adiabatic

γ_{HOT}

as function of
 q_0 and T_{HOT}
on-axis heating



Blue regions:
stabilizing

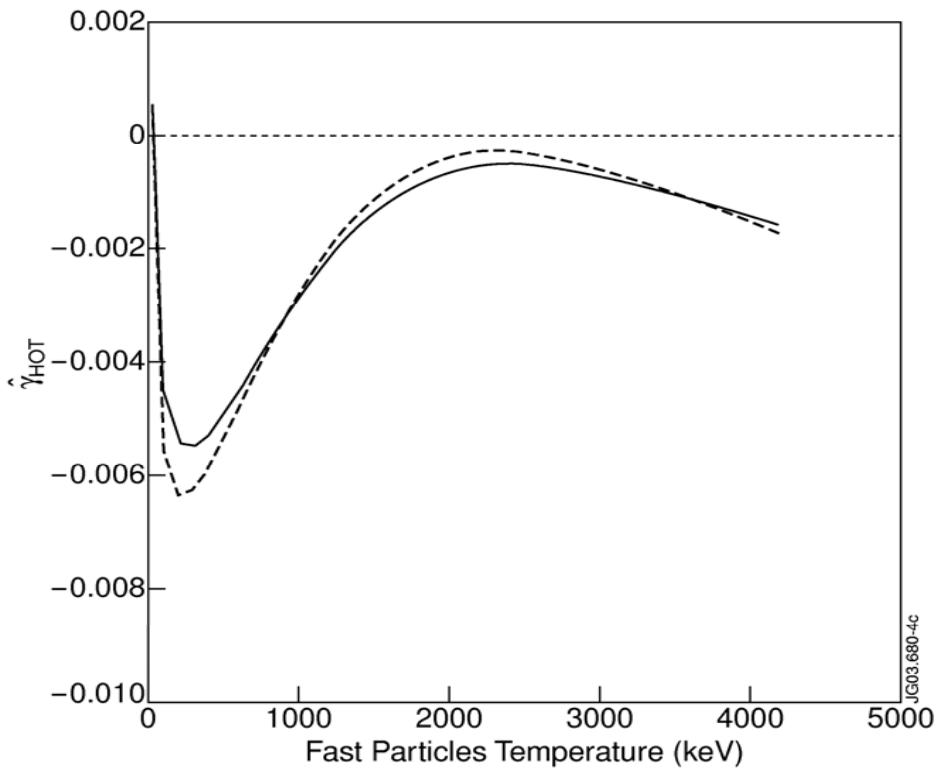
Red regions:
destabilizing

Adiabatic part of δW_{HOT}
(Preliminary results)

$\left\{ \begin{array}{l} \text{Destabilizing} \\ \delta W_{HOT}^{ad} \text{ increases with } T_{HOT} \end{array} \right.$

Application to experiments

→ Experiments with low plasma densities



Sawteeth destabilization
occurred when:

- ICRH power increases
- Plasma density decreases

$$\gamma_{HOT} \sim \frac{P_{ICRH}}{n_e}$$

Non-adiabatic γ_{HOT} as function of T_{HOT}

Application to experiments

- Sawteeth destabilization coincides with an increase in T_{HOT}

Numerical results show:

- The stabilizing effect of the non-adiabatic part of δW_{HOT} decreases as T_{HOT} increases, so this mechanism for sawtooth stabilization is weakened.
- The destabilizing effect associated with the adiabatic part of δW_{HOT} is increased.
- The non-adiabatic part of δW_{HOT} is dominant.

Variational method

$$\delta W_{MHD} + \delta W_{HOT} - \frac{8i\Gamma[(\Lambda^{3/2} + 5)/4][\omega(\omega - \omega_{*i})]^{1/2}}{\Lambda^{9/4}\Gamma[(\Lambda^{3/2} - 1)/4]\omega_A} = 0$$

(R. White et al.
Phys. Fluids B 2,
745 (1990))

Solutions of the internal kink dispersion relation: (Ideal limit)

	Kink Branch
	Sawteeth
Low frequency	
$\omega \approx \omega_{*i}$	
	Ion Branch
	Low frequency (diamagnetic)
High frequency	
$\omega \approx \langle \omega_{Dhi} \rangle$	
	Fishbone Branch
	High frequency (precessional)

(B. Coppi and F. Porcelli P. R.L, 57, 2272 (1986))

(L. Chen et. al P. R.L 52, 1122 (1984))

Fishbone Branch

High frequency (precessional) fishbones

Determination of the regions of stability

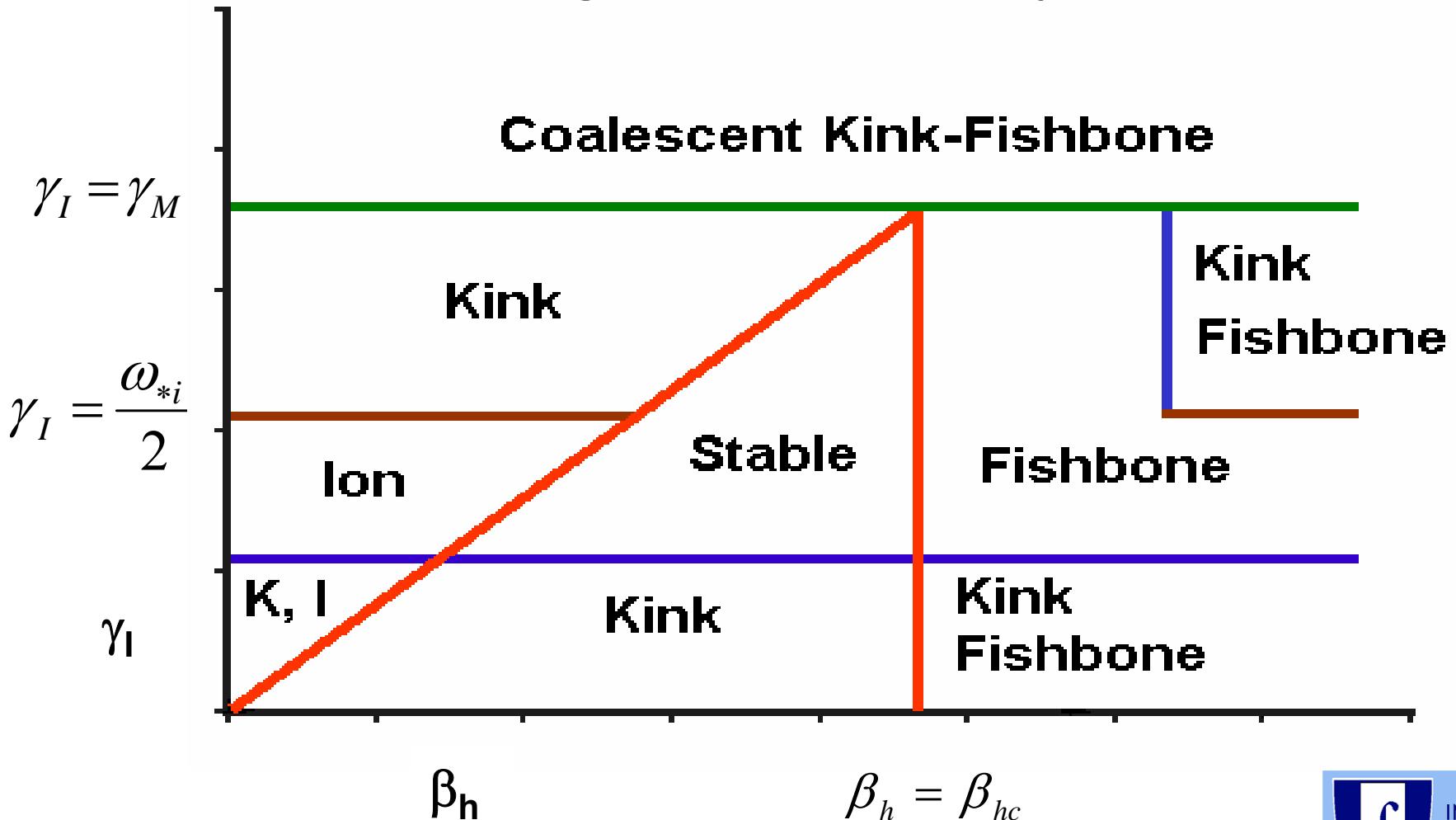
Stability governed mainly by	γ_I	Ideal (MHD) growth rate	$\gamma_I \equiv -\omega_A \delta W_{MHD}$
	ω_{*i}	Diamagnetic frequency	
	β_h	Fast particles beta	

Marginal equation $\omega = \text{real}$, ideal limit, ICRH population

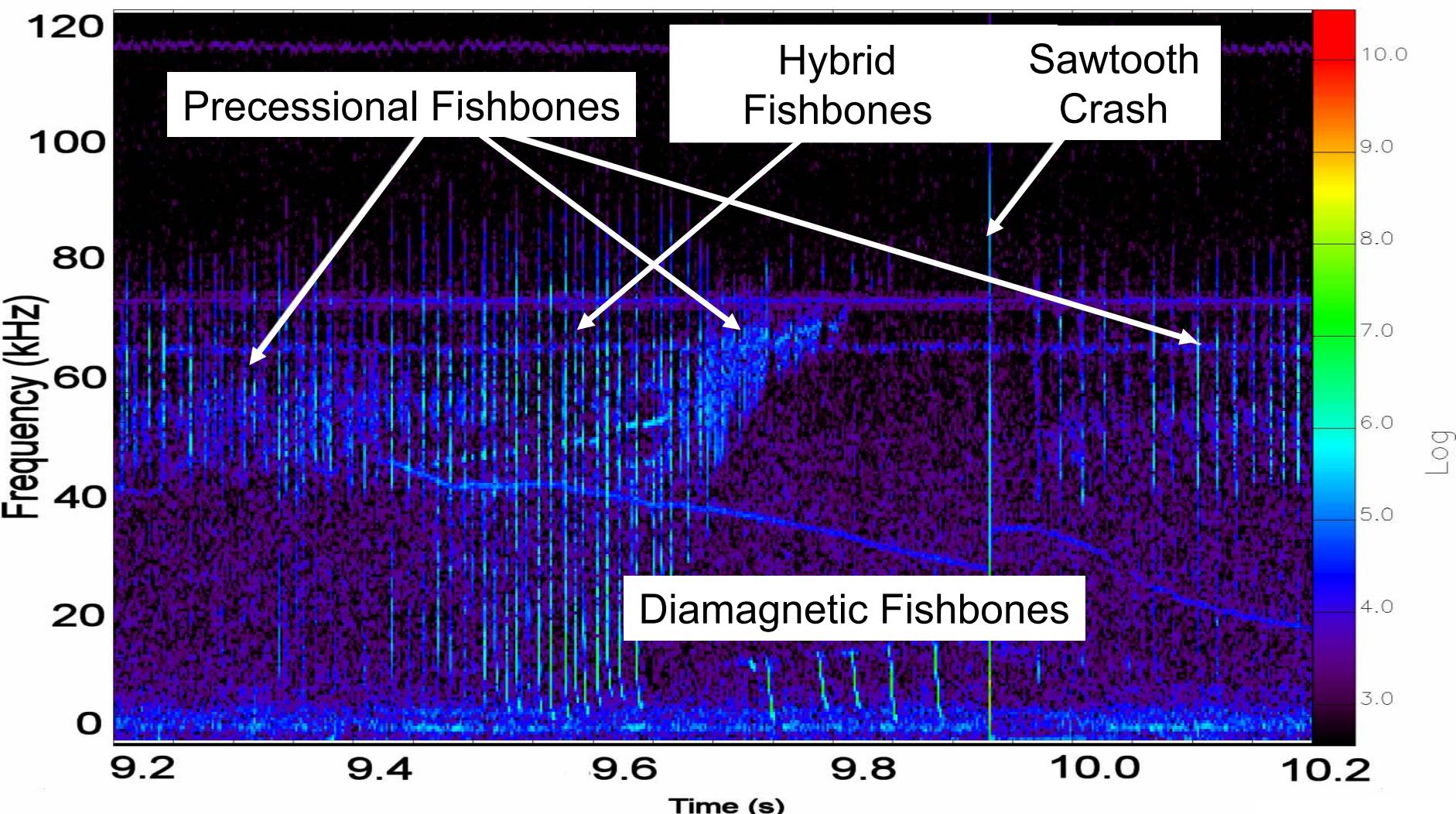
$$\underline{\gamma_I} = \frac{3}{4} \left[\frac{\omega}{\langle \omega_D \rangle} \left(\frac{\omega}{\langle \omega_D \rangle} - \frac{\omega_{*i}}{\langle \omega_D \rangle} \right) \right]^{\frac{1}{2}} \left(\frac{\omega}{\langle \omega_D \rangle} \right)^{-\frac{3}{2}} \left[\frac{1}{2} + \frac{\omega}{\langle \omega_D \rangle} + \left(\frac{\omega}{\langle \omega_D \rangle} \right)^{\frac{3}{2}} \text{Re} Z \left[\left(\frac{\omega}{\langle \omega_D \rangle} \right)^{\frac{1}{2}} \right] \right]$$

$$\underline{\beta_h} = \frac{3}{4} \frac{\varepsilon \omega_A}{\pi^{1/2} \langle \omega_D \rangle} \left[\frac{\omega}{\langle \omega_D \rangle} \left(\frac{\omega}{\langle \omega_D \rangle} - \frac{\omega_{*i}}{\langle \omega_D \rangle} \right) \right]^{\frac{1}{2}} e^{\omega / \omega_D} \left(\frac{\omega}{\langle \omega_D \rangle} \right)^{-\frac{5}{2}}$$

Regions of stability



Magnetic spectrogram from JET



Variation of the relevant parameters

γ_I 

Increases

In between crashes the
 $q=1$ surface expands

ω_{*i} 

Increases

In between crashes the
ion bulk profile peaks

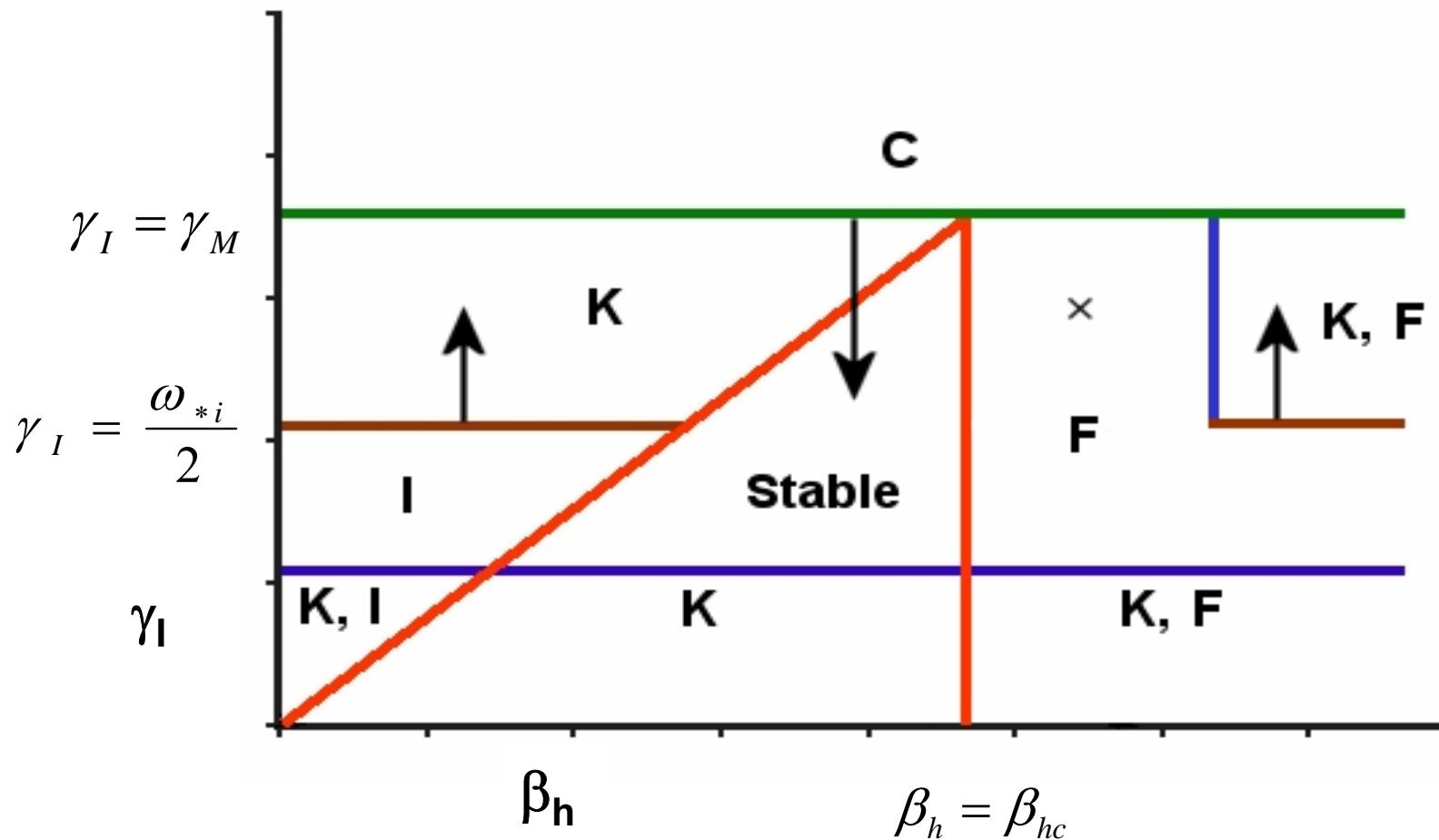
f_{*i} : 3 kHz  20 kHz

β_h 

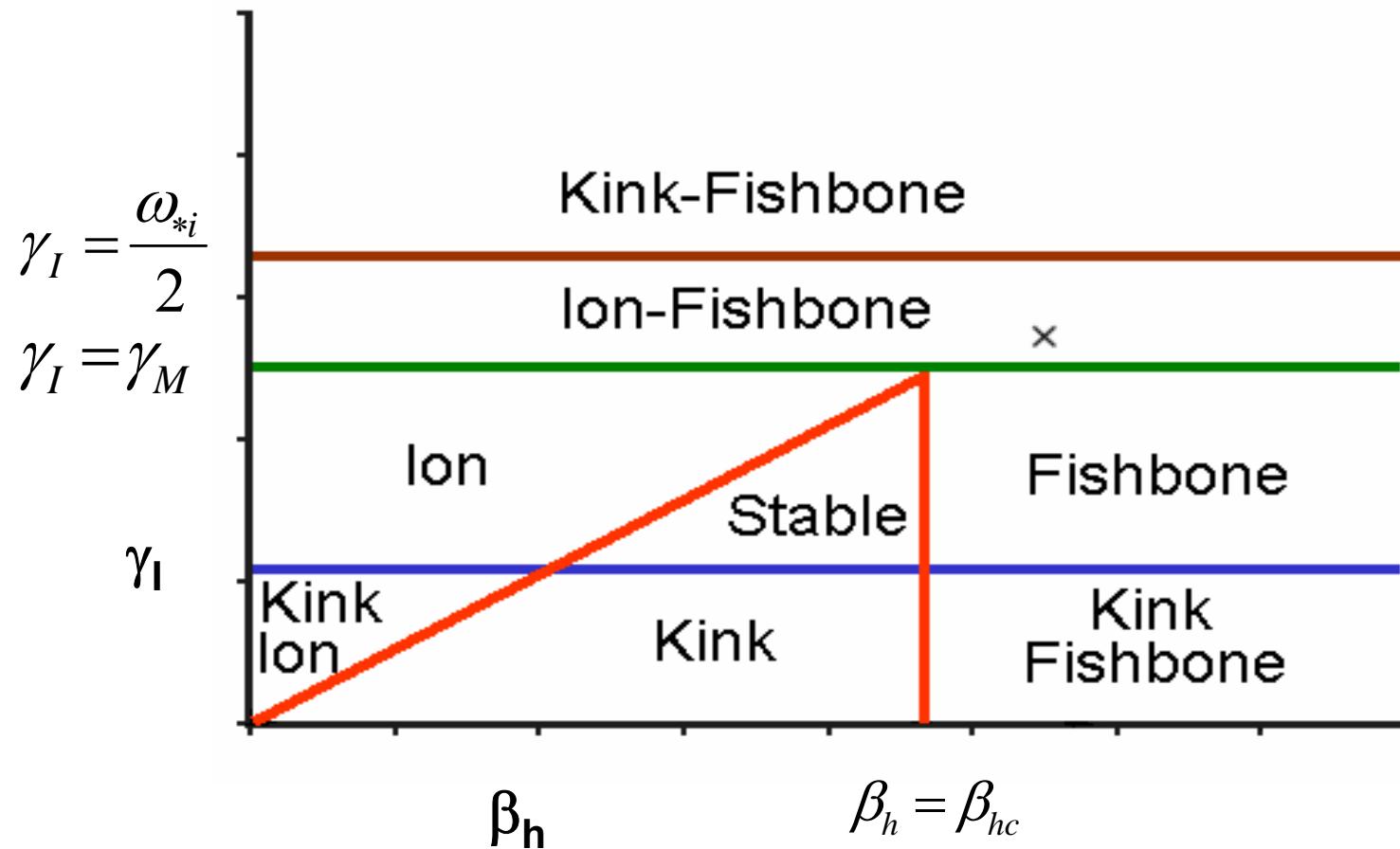
Decreases during a burst

Increases between bursts

Variation of the regions of stability

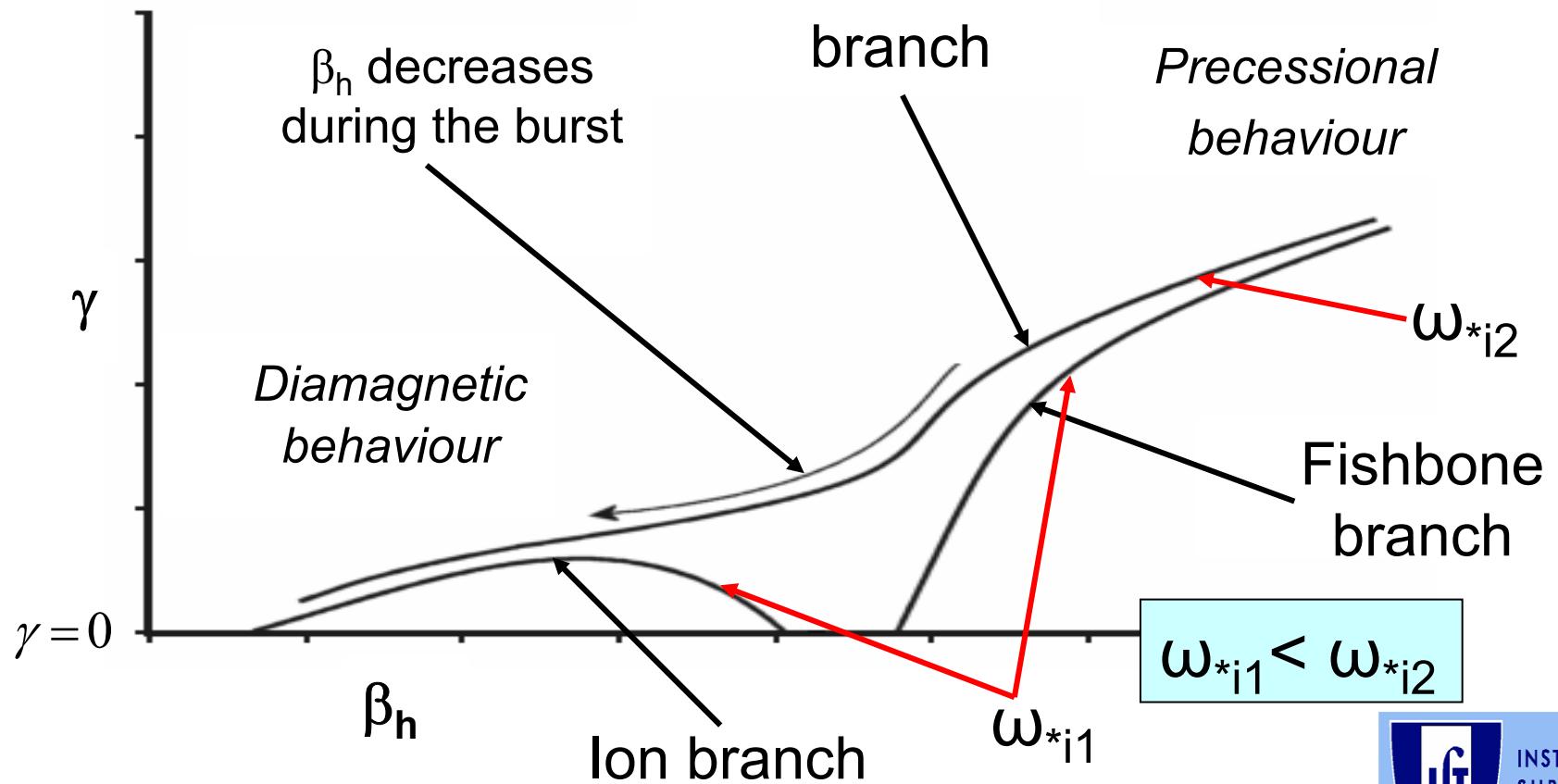


Regions of stability

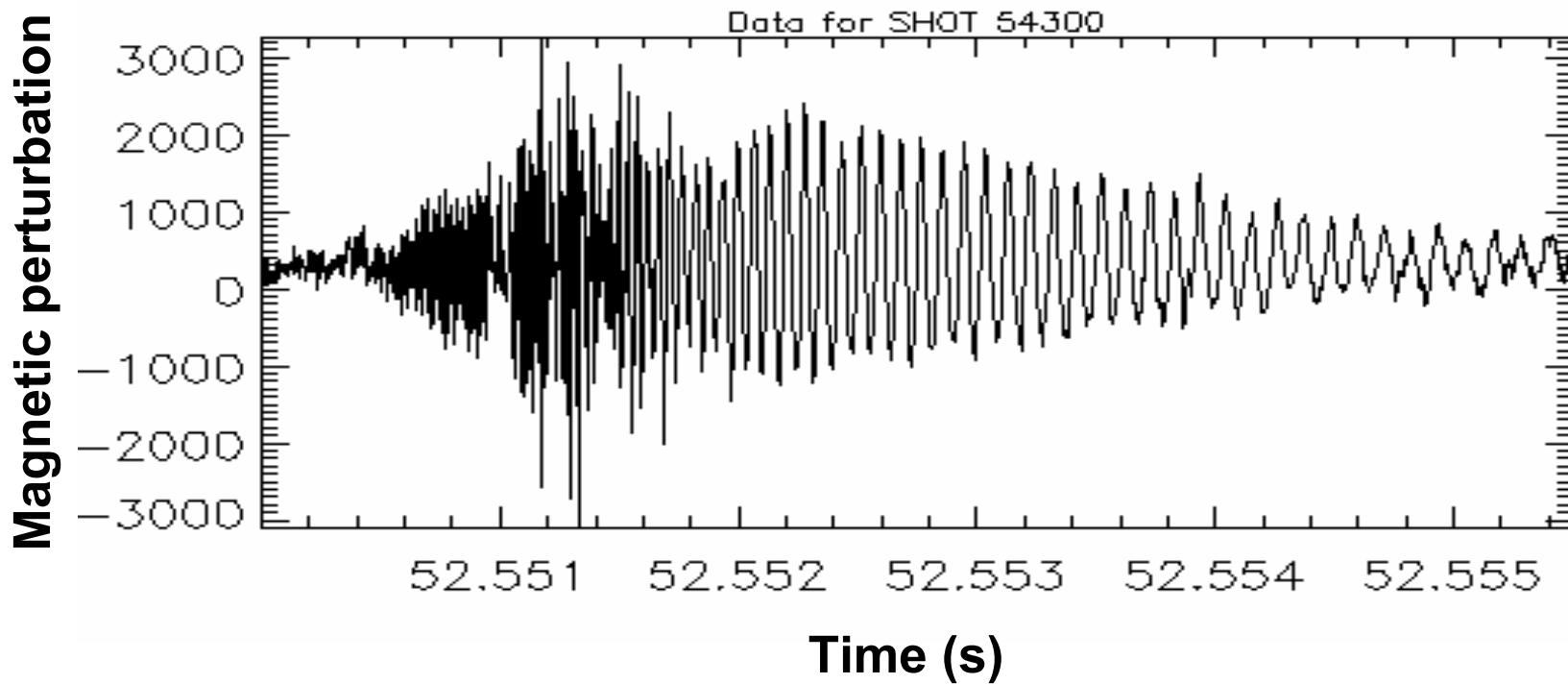


Solutions of the dispersion relation

(R. White et al. Phys. Fluids B 2, 745 (1990))



Magnetic perturbation for an hybrid fishbone



Modes and regimes

Ion mode



Low frequency
fishbones

Fishbone mode



High frequency
fishbones

(Gorolenkov *et al.* "Fast ions effects on fishbones and n=1 kinks in JET simulated by a non-perturbative NOVA-KN code, this conference)

Coalescent
ion-fishbone
mode



Hybrid fishbones

Both types of
fishbones occurring
simultaneously

Summary

Perturbative Method

- Accurate calculation of γ_{HOT} for a realistic geometry
- Only for the “kink mode”



Stabilizing term vanishes for high fast ions temperatures

Variational Method

- All branches of the dispersion relation
- Predicts changes in instabilities behaviour knowing the relevant parameters
- Simplified geometry and fast ions' distribution function, suitable only for a qualitative approach



Hybrid fishbones and both types of fishbones occurring simultaneously